

Nonoscillatory Central Schemes on Unstructured Triangulations for Hyperbolic Systems of Conservation Laws

Ivan Christov Bojan Popov

Department of Mathematics,
Texas A&M University, College Station, Texas 77843-3368

christov@tamu.edu

Minisymposium on Numerical Methods for First-Order PDEs,

8th IMACS International Symposium on
Iterative Methods in Scientific Computation,

College Station, Texas
November 16, 2006



- Motivation:
 - How should we design “genuinely multidimensional” limiters and (central Godunov-type) schemes on unstructured grids?
 - Need for a simple & fast predictor algorithm for the time-dependent extension of L^1 -minimization FEM (J.-L. Guermond & B. Popov).
- Background:
 - High-resolution central schemes: Nessyahu & Tadmor (1992) → Jiang & Tadmor (1998) → Kurganov & Tadmor (2000) → Kurganov, Noelle, Petrova (2001) ...
 - Extension to unstructured grids: Arminjon *et al.* (1997).
 - Semi-discere central-upwind version: Kurganov & Petrova (2005).
- Outline of the talk:
 - 1 Hyperbolic systems of conservation laws and 2D central schemes.
 - 2 The minimum-angle-plane reconstruction.
 - 3 Construction of the staggered grid corresponding to an *unstructured* triangulation.
 - 4 Numerical results for equations with *convex* and *nonconvex* fluxes.



Statement of the Problem and Notation

- Consider the initial-boundary-value problem for a 2D *hyperbolic* system of conservation laws:

$$\begin{cases} \vec{q}_t + \vec{f}(\vec{q})_x + \vec{g}(\vec{q})_y = 0, & (x, y, t) \in \Omega \times (0, T], \\ \vec{q}(x, y, t = 0) = \vec{q}_0(x, y), & (x, y) \in \Omega, \\ \vec{q}(x, y, t) = \vec{q}_{\text{BC}}(x, y, t), & (x, y, t) \in \partial\Omega \times (0, T]. \end{cases}$$

- $\Omega \subset \mathbb{R}^2$ is the interior of a polygonal domain and $\partial\Omega$ its boundary.
- $\mathcal{T} = \{\tau_i\}$ is a *conforming* triangulation of $\bar{\Omega}$.
- \bar{w}^n is a piecewise-constant approximation to the cell averages of \vec{q} on \mathcal{T} at time t^n .
- $\mathcal{S} = \{\sigma_k\}$ is the *staggered* grid — the “dual” of \mathcal{T} .
- \bar{w}^n is the analogue of \bar{w}^n on \mathcal{S} .



Overview of the 2D Central Scheme

1. Perform a slope-limited piecewise-linear reconstruction on \mathcal{T} .

$$\bar{w}^n \longrightarrow w^n$$

2. Evolve the cell averages on the staggered grid \mathcal{S} in time.

$$w^n \longrightarrow \bar{w}^n, \quad \bar{w}^n \longrightarrow \bar{w}^{n+1}$$

3. Project the solution from \mathcal{S} back onto \mathcal{T} .

$$\bar{w}^{n+1} \longrightarrow w^{n+1}, \quad w^{n+1} \longrightarrow \bar{w}^{n+1}$$

The good, the bad and the ugly:

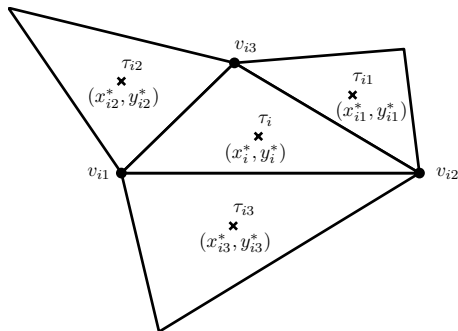
- No need to solve a Riemann problem at each cell interface!
- Need to define \mathcal{S} in a “reasonable” manner.
- Need to be able to perform a nonoscillatory reconstruction on \mathcal{S} .



The Minimum-Angle-Plane Reconstruction

The algorithm:

- 1 Given an element $\tau_i \in \mathcal{T}$ and its neighbors τ_{ij} , $1 \leq j \leq m$, find all $\binom{m+1}{3}$ possible planes.
- 2 Find the plane that makes the smallest angle with the horizontal, and use it to find a *limited* gradient.



Note that

- This “genuinely 2D” limiter behaves like minmod with a UNO flavor (i.e., \approx Durlofsky–Engquist–Osher but $> 1^{\text{st}}$ order near extrema).
- No particular geometry and/or connectivity is assumed in the design of the limiter, and there are no *ad hoc* parameters.
- Limiter works exactly the same way on \mathcal{S} as on \mathcal{T} .



The Staggered Grid

The dual elements:

- 1 Triangles Δ_j .
- 2 Polygons Λ_{ij} .
- 3 Parallelograms Π_{ij} .

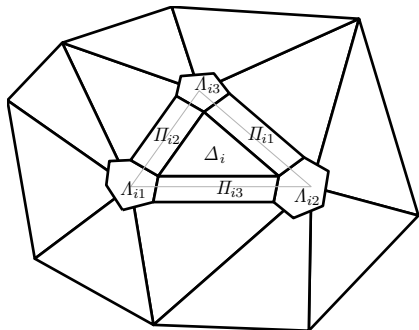
The usual CFL condition

$$\Delta t < \frac{1}{3} \cdot \min_i |\tau_i| / S_{\max},$$

S_{\max} = fastest wave's speed,
is good enough.

In particular:

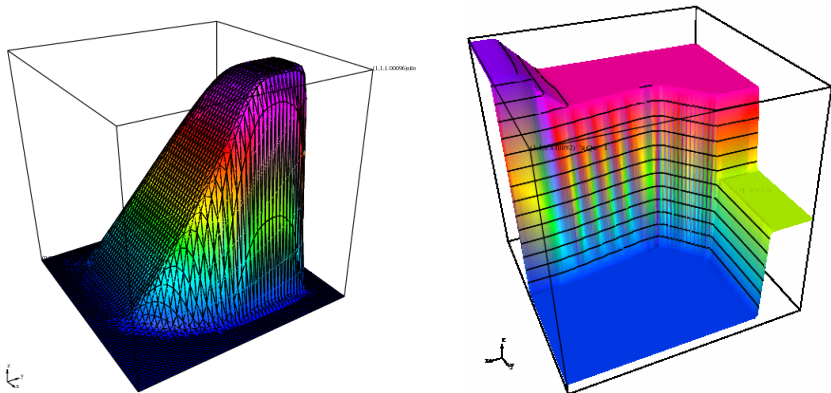
- If $|\Delta_i| = |\Pi_{ij}| = 0$, the staggered grid becomes the Voronoi diagram.
- If the maximum local speed of propagation is used, we get the analogue of Kurganov & Tadmor's modified central differencing for triangulations.



Numerical Results for a *Convex Flux*

- Riemann problems for the 2D inviscid Burgers equation

$$u_t + \left(\frac{1}{2}u^2\right)_x + \left(\frac{1}{2}u^2\right)_y = 0.$$

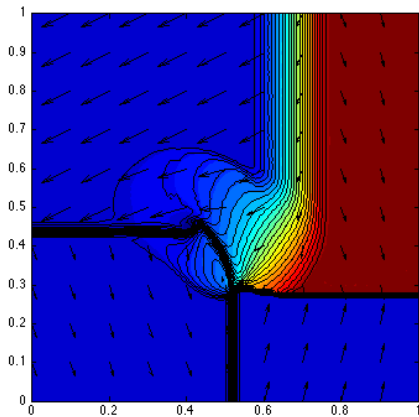
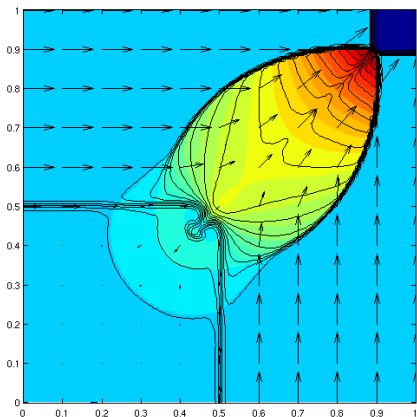


- (L) 6,272 elements and 19,041 dual elements.
- (R) 12,800 elements and 38,721 dual elements.



Riemann Problem for the Euler Equations

- $\vec{q} = (\rho, \rho u, \rho v, E)^\top$, $\vec{f}(\vec{q}) = (\rho u, \rho u^2 + p, \rho uv, u(E + p))^\top$,
 $\vec{g}(\vec{q}) = (\rho v, \rho uv, \rho v^2 + p, v(E + p))^\top$,
- For an ideal gas: $p = (\gamma - 1) [E - \frac{1}{2}\rho(u^2 + v^2)]$, $\gamma = 1.4$.



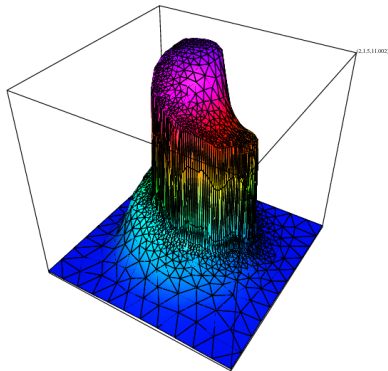
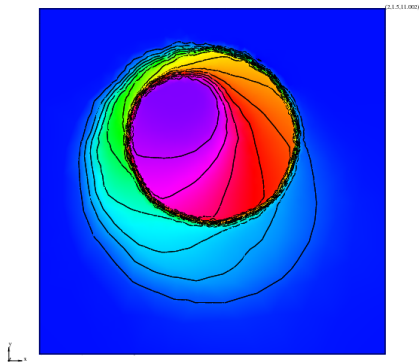
- Very fine mesh ($\min_i \text{diam}(\tau_i) = \sqrt{2}/256$) with CFL = 0.275.



Numerical Results for a *Nonconvex* Flux (I)

- Riemann problem for the 2D scalar equation

$$u_t + (\sin u)_x + (\cos u)_y = 0.$$

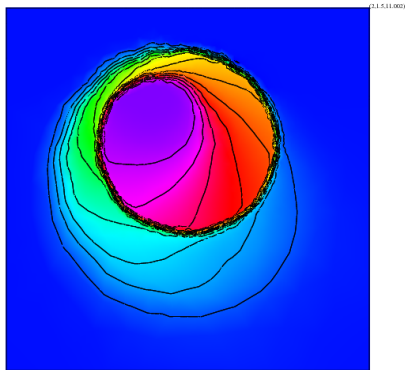
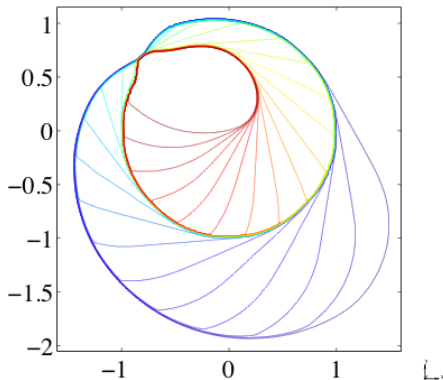


- Adapted mesh with (only) 3,264 elements and 9,837 dual elements, $\text{CFL} = \frac{1}{6}$.



Numerical Results for a *Nonconvex* Flux (II)

- Kurganov, Petrova & B. Popov reported that less compressive / higher order limiters (e.g., WENO5, MM2, SB) do *not* resolve the resulting composite wave correctly. **The MAPR passes this test!**



- (L) **90,000** element Cartesian tensor-product mesh.



Selected Bibliography

-  D. Kröner, *Numerical Schemes for Conservation Laws*. Wiley–Teubner, 1997.
-  G.-S. Jiang & E. Tadmor, “Nonoscillatory Central Schemes for Multidimensional ... ”
SIAM J. Sci. Comput. **19** (1998) 1892.
-  A. Kurganov & G. Petrova, “Central-Upwind Schemes on Triangular Grids ... ”
Numer. Methods Partial Differential Eq. **21** (2005) 536.
-  A. Kurganov, G. Petrova & B. Popov, “Adaptive Semi-Discrete Central-Upwind Schemes ... ”
SIAM J. Sci. Comput. (submitted).
-  J.-L. Guermond, “A Finite Element Technique for Solving First-Order PDEs in L^p .”
SIAM J. Numer. Anal. **42** (2004) 714.

